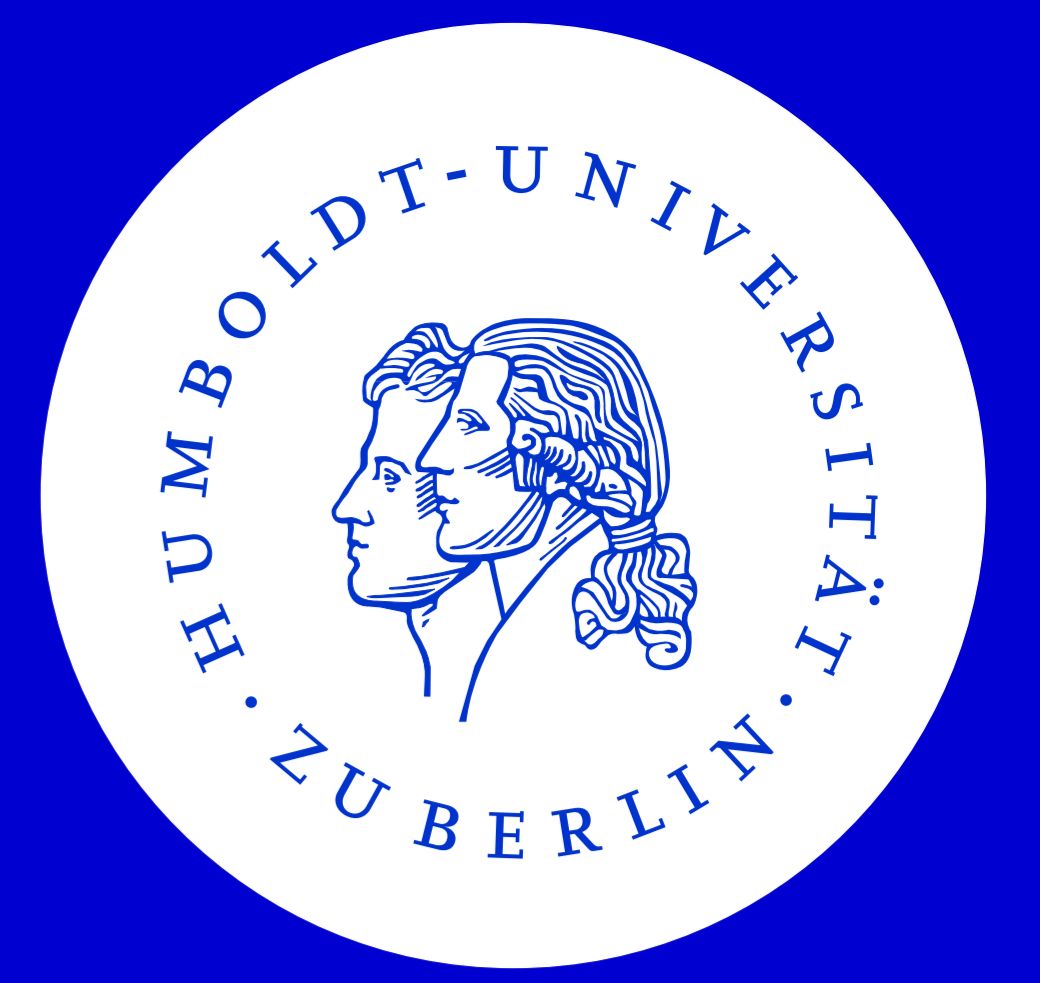


# INELASTIC RESONANCES IN 1D AND 2D CONFINEMENT

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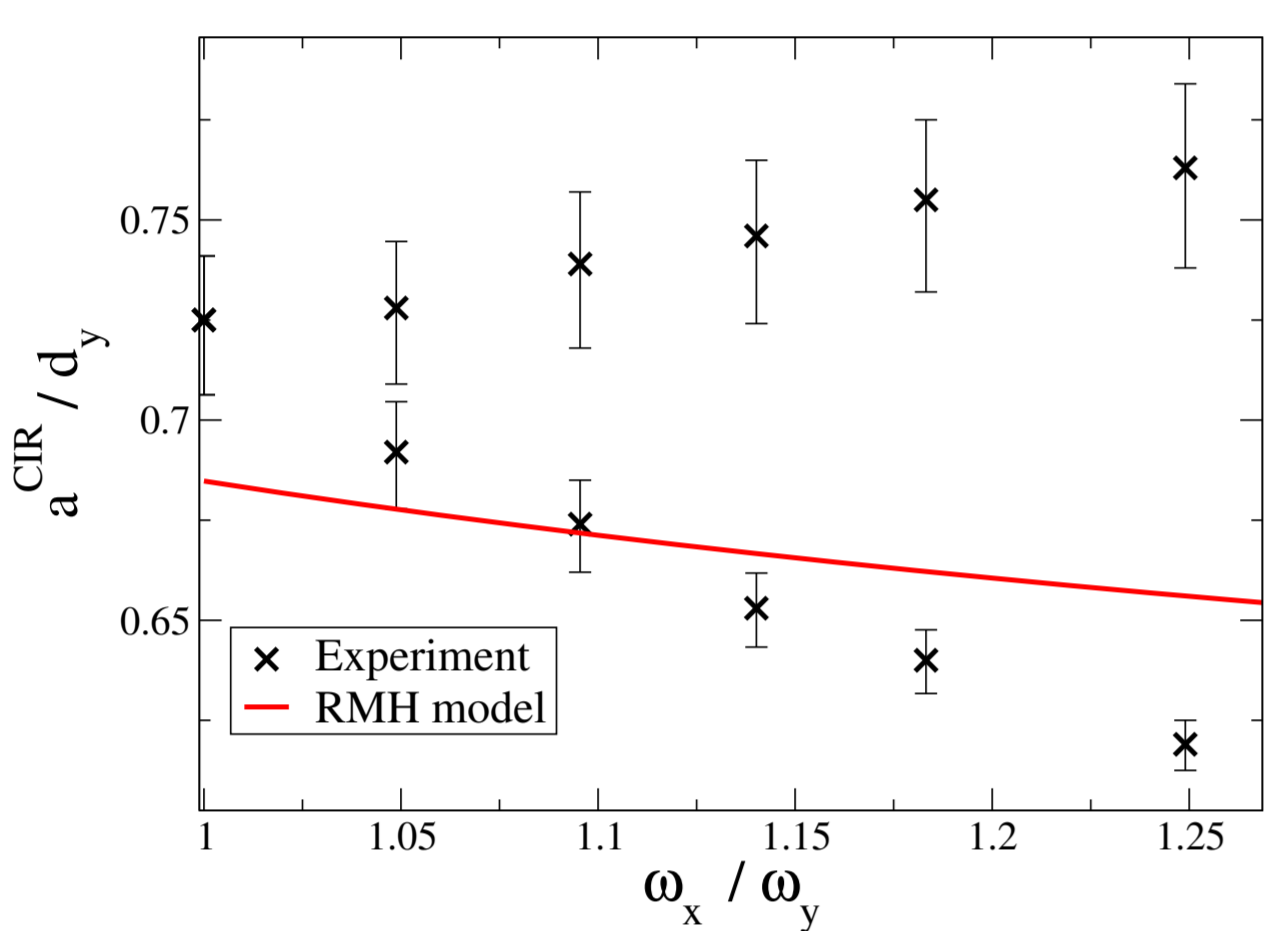


## Motivation

The established theories of **confinement-induced resonances** (CIRs) [1,2,3] predict divergences of the effective interaction strength  $g$  for strong anisotropic confinement. These models are based on the relative motion Hamiltonian of two atoms in an harmonic trapping potential (RMH model). Experimentally, the divergence of  $g_{2D}$  was very recently demonstrated using rf spectroscopy [4]. However, an analysis of the RMH models [5,6] shows that they seem to contradict recent experimental observations of loss resonances in strongly anisotropic traps [7]. In this work [8], we present a model based on the coupling of relative and center-of-mass motion (CRC model) due to the anharmonicity of the external potential which yields (true) Feshbach resonances. We compare the resulting inelastic resonances to the loss resonances in [7].

## Elastic vs. Loss Resonances in 1D

1D confinement is characterized by  $\omega_x, \omega_y \gg \omega_z$ . The divergence of the effective interaction strength  $g_{1D}$  predicted by the RMH model is an elastic scattering process which may not be visible in the loss experiment of Haller *et al.* [7]:



In the plot a comparison of the **elastic CIR** positions predicted by the RMH model (red line) with the **loss resonances** observed in the experiment [7] for different values of the anisotropy of the transversal confinement is shown. While the loss resonances show a **splitting** for transversal anisotropy, the elastic CIR does not split as proven in [5,6]. Hence, **the elastic RMH model is not capable of describing the loss resonances of the experiment [7]**.

⇒ Different behavior of elastic CIR and loss resonances in 1D.

## Elastic vs. Loss Resonance in 2D

2D confinement is characterized by  $\omega_x \gg \omega_y, \omega_z$ . As in 1D, an elastic CIR is predicted by the divergence of the effective interaction strength  $g_{2D}$  by an RMH model [2]. This **elastic CIR occurs only for negative values of the scattering length** and was recently confirmed experimentally [4]. However, the experiment [7] observes a 2D loss resonance only at a positive value of the scattering length.

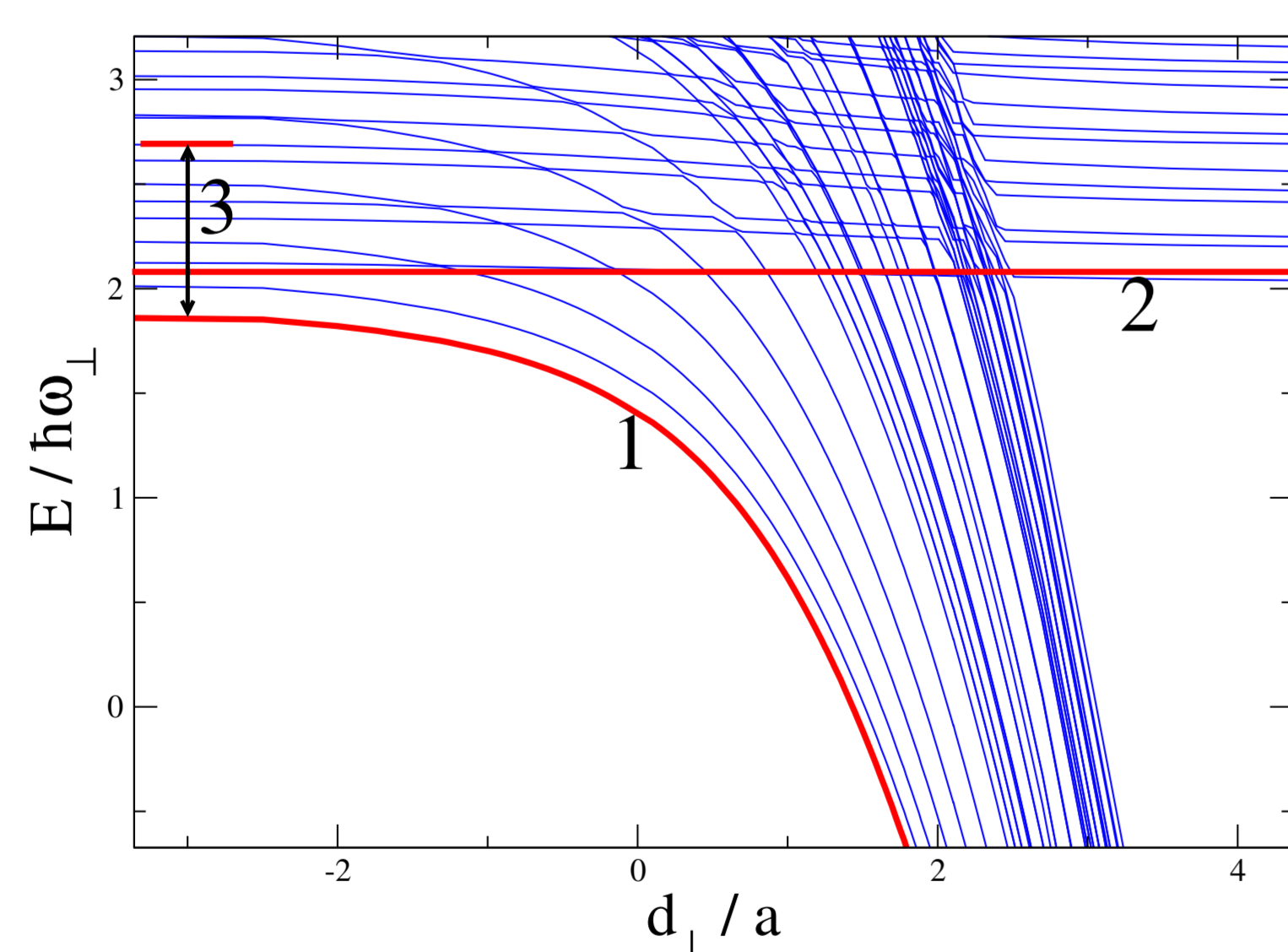
⇒ Different positions of elastic CIR and loss resonance in 2D.

## Universal Model for Inelastic CIRs

### Idea:

In the experiment [7] an optical lattice is used leading to a coupling of relative (REL) and center-of-mass (COM) motion. The COM-REL coupling leads to **Feshbach resonances** induced by (avoided) crossings of bound states with COM excitation and a state of an unbound atom pair without COM excitation. At the avoided crossing a molecular bound state can be occupied, since the excess binding energy can be transferred to COM excitation energy, which is a **inelastic** process. The **formation of molecules** is a loss mechanism and should be observable in loss experiments like [7]. The proposed model is **universally** valid for 1D as well as 2D confinement.

### Energy Formula:

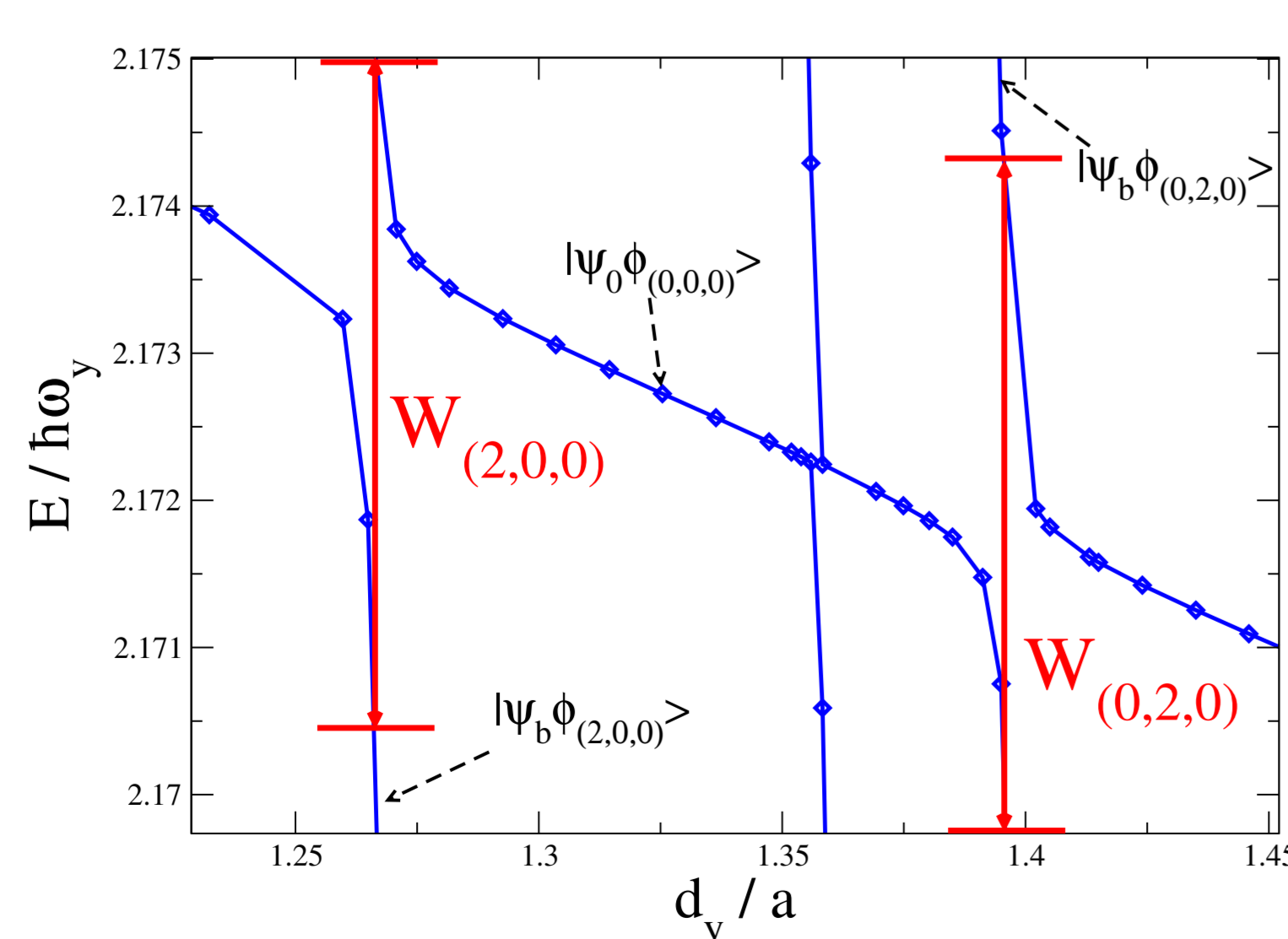


- 1 Bound state energy  $E_b^{\text{REL}}(a)$  (using harmonic approximation)
- 2 Energy  $E_1$  of the lowest trap state: (using constant approximation)
- 3 Center-of-mass excitation  $\Delta_n$  (using 1st-order perturbation theory)

Crossings occur for:

$$E_b^{\text{REL}} = E_1^{\text{REL}} - \Delta_n$$

### Matrix Elements and CIR Selection Rules in 1D:



$$W_n \approx \delta_{n_z,0} \left[ \delta_{n_y,0} \langle \phi_{n_x} \psi^{(b)} | W_x | \psi_1 \phi_0 \rangle + \delta_{n_x,0} \langle \phi_{n_y} \psi^{(b)} | W_y | \psi_1 \phi_0 \rangle \right]$$

- Resonances only for **transversally** (tight direction:  $\omega_x$ , or  $\omega_y$ ) excited bound states.
- $\mathbf{n} = (2, 0, 0)$  and  $\mathbf{n} = (0, 2, 0)$  are dominant.  $\mathbf{n} = (n_x, n_y, n_z)$
- Degenerate for  $\omega_x = \omega_y$ , **splitting** for  $\omega_x \neq \omega_y$

### Matrix Elements and CIR Selection Rules in 2D:

$$W_n \approx \delta_{n_y,0} \delta_{n_z,0} \langle \phi_{n_x} \psi^{(b)} | W_x | \psi_1 \phi_0 \rangle$$

- Resonances only for **transversally** (tight direction:  $\omega_x$ ) excited bound states.
- $\mathbf{n} = (2, 0, 0)$  is dominant.

## Ab Initio Calculation

### Hamiltonian (2 atoms in optical lattice (OL)):

$$H(\mathbf{r}, \mathbf{R}) = H_{\text{REL}}(\mathbf{r}) + H_{\text{COM}}(\mathbf{R}) + W(\mathbf{r}, \mathbf{R})$$

$$H_{\text{REL}}(\mathbf{r}) = T_{\text{REL}}(\mathbf{r}) + V_{\text{REL}}(\mathbf{r}) + U_{\text{int}}(r)$$

$$H_{\text{COM}}(\mathbf{R}) = T_{\text{COM}}(\mathbf{R}) + V_{\text{COM}}(\mathbf{R})$$

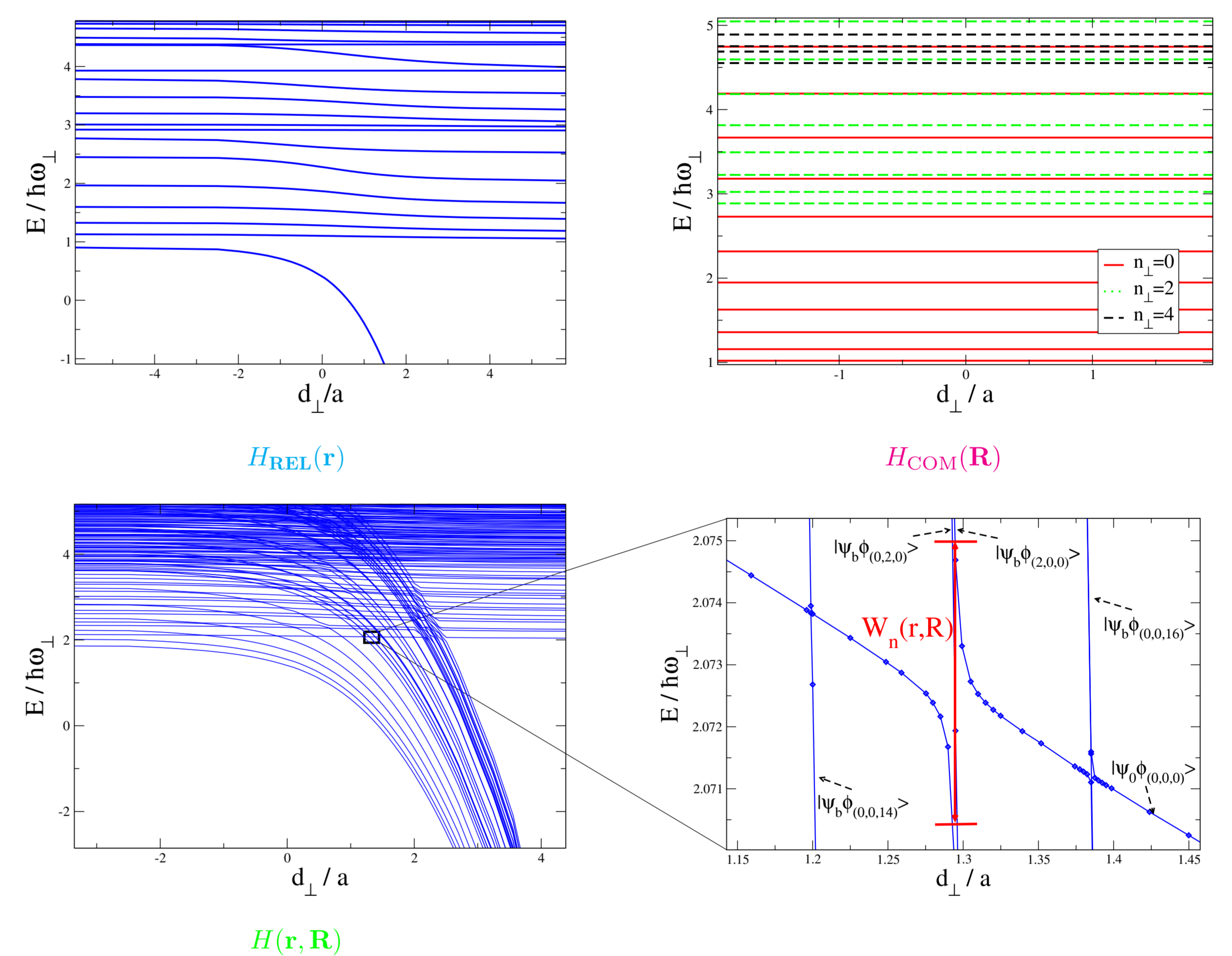
$\mathbf{r} = \mathbf{r}_1 - \mathbf{r}_2$  relative motion (REL) coordinate  
 $\mathbf{R} = \frac{1}{2}(\mathbf{r}_1 + \mathbf{r}_2)$  center of mass motion (COM) coordinate  
 $T_{\text{REL}}(\mathbf{r}), T_{\text{COM}}(\mathbf{R})$  kin. energy operators  
 $V_{\text{REL}}(\mathbf{r}), V_{\text{COM}}(\mathbf{R})$  separable parts of OL potential  
 $W(\mathbf{r}, \mathbf{R})$  non-separable OL terms  
 $U_{\text{int}}(r)$  interatomic interaction potential

### Exact Diagonalization:

The energy spectrum of the full Hamiltonian is calculated. To include the COM-REL coupling the optical lattice (OL) potential is expanded in a Taylor series and truncated after the 6th power (sextic potential). The interatomic interaction is included using full Born-Oppenheimer (BO) potential curves. The scattering length is varied by an inner-wall manipulation of the BO curves.

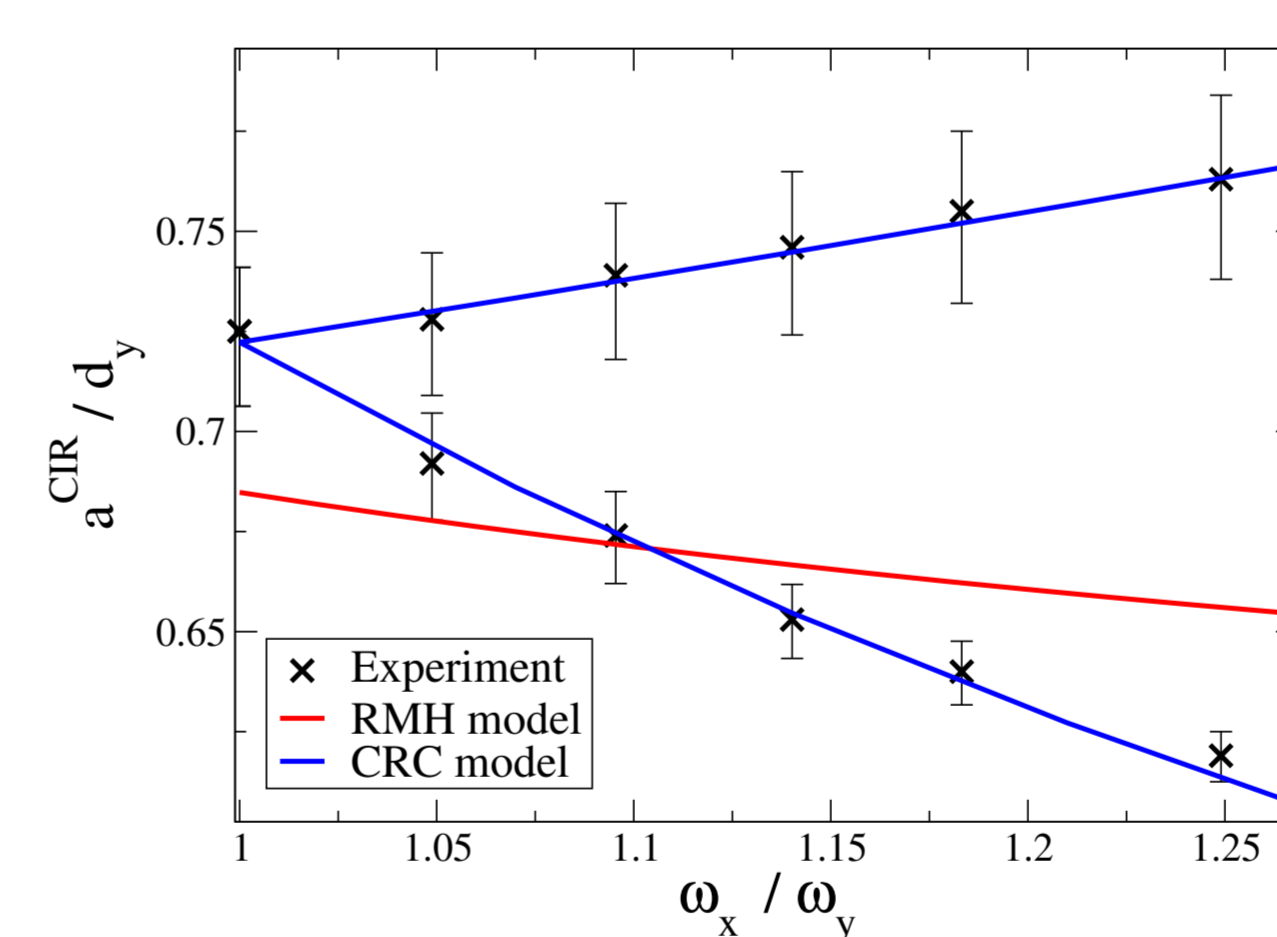
### Composition of Energy Spectra:

Energy spectra for two atoms confined in a sextic trap with  $\omega_x/\omega_z = \omega_y/\omega_z = 10$ .



## Comparison with Experiment

### CIR Positions in Quasi 1D:



### CIR Positions in Quasi 2D:

Method	$a^{(\text{CIR})}/d_y$
Experiment	0.593
CRC model	0.595
RMH model	negative value

Perfect agreement of inelastic CIR with loss resonances of experiment.

## Conclusion

Coupling between center-of-mass and relative motion leads to **inelastic confinement-induced resonances**. The experiment of Haller *et al.* [7] measured **inelastic losses** rather than directly the effective interaction strength. Hence, it is much more sensitive to the here present inelastic CIRs than to the elastic ones predicted by RMH models [1,2]. As **perfect agreement** of the presented model with experiment [7] shows, **the measured loss resonances are inelastic CIR which occur because of COM-REL coupling due to the anharmonicity of the external confinement**. (For more details see arXiv:1104.1561 [8].)

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